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An Adjoint Gamma-Ray Moments Computer Code, ADJMOM-I

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An Adjoint Gamma-Ray Moments Computer Code, ADJMOM-I

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 Washington, D.C. 20234

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George L. Simmons**

In this paper we discuss a computer code for generating spatial-angular moments of the adjoint gamma-ray flux in an infinite medium. The equation for the flux moments is given and techniques used for the solution are described. Details of the input data and a sample problem are also supplied.

Key words: Adjoint; buildup factor; dosimetry; gamma-ray transport; moments method; shielding.

1. Introduction and Theory

This note describes a computer code, ADJMOM-I, which can be used to calculate the moments (spatial and angular) of the adjoint gamma-ray flux distribution in an infinite homogeneous medium. These moments can then be used to calculate dose distributions for monoenergetic point isotropic, plane isotropic, plane oblique, and point conical sources. The adjoint moments method is simply the application of the moments method [1]¹ to the adjoint gamma-ray transport equation. A particular feature of the adjoint formulation is that the contribution of annihilating electron pairs and fluorescent gamma-rays to the gamma-ray transport is quite simple to include. Additionally the adjoint source may have delta function angular characteristics thereby allowing the calculation of the dose angular distribution.

The adjoint moments method for gamma-rays has been applied by several authors [2,3,4,5] but only a limited amount of results have been

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¹ Figures in brackets indicate the literature references at the end of this paper.

obtained. In the formulation of the adjoint equation, we begin with the one dimensional gamma-ray transport equation for a plane monoenergetic source located at $z = 0$ with angular characteristics $S(\vec{\omega})$,

$$\begin{aligned} \cos\theta \frac{\partial N}{\partial z}(E, z, \vec{\omega}) + \mu(E)N(E, z, \vec{\omega}) &= \int_E^{E_0} dE' \int_{4\pi} \frac{d\Omega'}{2\pi} K(E, E', \cos\theta) N(E', z, \vec{\omega}) + \\ &+ S(\vec{\omega}) \frac{\delta(E-E_0)\delta(z)}{4\pi} + \frac{\delta(E-1)}{2\pi} \int_2^{E_0} dE' \int_{4\pi} d\Omega' \mu_{pp}(E') N(E', z, \vec{\omega}') + \\ &+ \omega_k \frac{\delta(E-E_k)}{4\pi} \int_{E'_k}^{E_0} dE' \int d\Omega' \mu_{ph}(E') N(E', z, \vec{\omega}'). \end{aligned} \quad (1)$$

where the last two terms in the equation represent the contribution of annihilating electron pairs and fluorescent gamma-rays, respectively.

Note that electron motion is not considered. In Equation (1)

$N(E, z, \vec{\omega})$ is the photon flux (gammas/sr-sec-cm²),

E is the photon energy ($m_0 c^2$ units),

$\vec{\omega}$ is the photon direction with $\vec{\omega} \cdot \vec{k} = \cos\theta$, when \vec{k} is a unit vector in the positive z direction.

E_0 is the source energy,

$\mu(E)$ is the total attenuation coefficient in Thomson units per electron (TU/el), $\mu(\text{cm}^2/\text{g}) \approx 0.400594 (Z/A) \mu(\text{TU/el})$,

$\cos\theta = 1 - \frac{1}{E} + \frac{1}{E'}$,

E' is the initial energy of the gamma ray,

E is the final energy after scattering,

$\mu_{pp}(E)$ is the pair production cross section (TU/electron),

$\mu_{ph}(E)$ is the photoelectric absorption cross section (TU/el),

E_k is the fluorescent yield for K X-rays of energy E_k ,

E'_k is the minimum energy for ionizing K electrons,

$$K(E, E') = 0.375 E \left\{ \frac{E}{E'} + \frac{E'}{E} + \cos^2 \Theta - 1 \right\} / (E')^3.$$

In developing the equation adjoint to Equation (1), i.e. the forward equation, we note that the forward equation may be simply written as

$H N = S$, where H is the forward transport operator, N is the flux and S is the source term. The adjoint equation can be represented as

$H^\dagger N^\dagger = S^\dagger$, where we use the dagger to indicate the adjoint analogues to the forward equation. If these two equations are to be adjoint, then in operation notation we require

$$(N, H^\dagger N^\dagger) \equiv (N^\dagger, HN). \quad (2a)$$

Clearly, a consequence of this relation is the requirement that

$$(N, S^\dagger) = (N^\dagger, S). \quad (2b)$$

It is this property of the adjoint flux which makes it useful in calculating dose distributions for many types of sources. That is to say, if we use the product of the energy deposition coefficient, $\mu_{ed}(E)$ [8], and photon energy, E , as the adjoint source, the energy deposition distribution for the forward problem may be obtained by performing the integrals indicated in Equation (2b). Particularly if

$$S(E, \vec{\omega}, z) = \frac{\delta(E - E_0) \delta(z)}{4\pi}$$

and

$$S^\dagger(E, \vec{\omega}, z) = \frac{E \mu_{ed}(E) \delta(z)}{4\pi}$$

then

$$\begin{aligned} D(z) &= \int_0^{E_0} dE' E' \mu_{ed}(E') \int_{4\pi} d\Omega' N(E', z, \vec{\omega}') \\ &= \int_0^{\infty} dE' \delta(E' - E_0) \int_{4\pi} d\Omega' N^\dagger(E', z, \vec{\omega}') \end{aligned} \quad (2c)$$

It is the property of the adjoint flux which allows the computation of dose distributions for many monoenergetic types.

The derivation of the adjoint transport equation is well known and will not be given here. The reader is referred to Reference [6] for a detailed discussion of the formulation of this equation as well as for an extensive bibliography on the subject. The equation for the adjoint flux for a detector with angular response $R(\vec{m})$ is

$$\begin{aligned} -\cos\theta \frac{\partial}{\partial z} N^\dagger(E, z, \vec{\omega}) + \mu(E) N^\dagger(E, z, \vec{\omega}) &= \int_0^E dE' \int \frac{d\Omega'}{2\pi} K(E', E, -\cos\theta) \times \\ &\times N^\dagger(E', z, \vec{\omega}') + E \mu_{ed}(E) \delta(z) R(\vec{\omega}) + \frac{\mu_{pp}(E)}{2\pi} \int_{4\pi} d\Omega' N^\dagger(1, z, \vec{\omega}') + \\ &+ \frac{\omega \mu_{ph}(E)}{2\pi} \int_{4\pi} d\Omega' N^\dagger(E_k, z, \vec{\omega}'). \end{aligned} \quad (3)$$

where E_c is the minimum final energy attained by a gamma ray of energy E undergoing a single Compton scattering.

Note that Equation (2c) may be obtained by applying Equation (2a) to Equations (1) and (3), with $R(\vec{m}) = S(\vec{m}) = 1$.

If we define adjoint moments as

$$N_{n,\ell,m}^{\dagger}(E) = \frac{\mu(E)^n}{4\pi E n!} \left(\frac{4\pi}{2\ell+1} \right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dz \, z^n \int_{4\pi} d\Omega' \, Y_{\ell}^{m*}(\theta, \varphi) N^{\dagger}(E, z, -\vec{\omega}), \quad (4)$$

then Equation (3) can be transformed into the following adjoint moments equation

$$\begin{aligned} \mu(E) N_{n,\ell,m}^{\dagger}(E) = & \int_{E_c}^E dE' \frac{E'}{E} K(E', E) P_{\ell} \left(1 + \frac{1}{E} - \frac{1}{E'} \right) \left[\frac{\mu(E)}{\mu(E')} \right]^n N_{n,\ell,m}^{\dagger}(E') + \\ & + \delta_{no} R_{\ell,m}^{\mu_{ed}}(E) + \frac{(1-\delta_{no})\mu(E)}{2\ell+1} \left\{ \left[(\ell+1)^2 - m^2 \right]^{\frac{1}{2}} N_{n-1,\ell+1,m}^{\dagger}(E) + \right. \\ & + \left. (\ell^2 - m^2)^{\frac{1}{2}} N_{n-1,\ell-1,m}^{\dagger}(E) \right\} + 2\delta_{\ell o} \frac{\mu_{pp}(E)}{E} \left[\frac{\mu(E)}{\mu(1)} \right]^n N_{n,o,o}^{\dagger}(1) + \\ & + \omega_k \delta_{\ell o} \frac{\mu_{ph}(E)}{E} \left[\frac{\mu(E)}{\mu(E_k)} \right]^n N_{n,o,o}^{\dagger}(E_k). \end{aligned} \quad (5)$$

where $R_{\ell,m}$ is the spherical harmonic coefficient of the angular response function $R(\vec{\omega})$.

The adjoint moments code ADJMOM-I calculates solutions to equation (5). Whereas in the forward moments solution, the solution procedure starts at $E = E_{\max}$ and proceeds downward to some minimum solution energy E_{\min} , in the adjoint moments solution, the solution procedure starts with E_{\min} (typically 10 keV) and goes upward in energy to some maximum energy E_{\max} (typically 15 MeV). The technique used in ADJMOM-I to solve Equation (5) is similar to the one that is used in the neutron moments code, MOMENT-I[7]. It involves the use of Gaussian Quadrature. We want to obtain solutions at N energies logarithmically spaced between

E_{\max} and E_{\min} ,

$$E_{\min} = E_1 < E_2 < E_3 < \dots < E_{N-1} < E_N = E_{\max},$$

where we require that $E_i > \frac{E_{i+1}}{2E_{i+1} + 1}$. To illustrate the use of Gaussian Quadrature, we consider Equation (5), for the $m = 0$ case and an isotropic detector:

$$\begin{aligned} \mu(E) N_{n,\ell}^{\dagger}(E) &= \int_{E_{i-1}}^E dE' \frac{E'}{E} K(E', E) P_{\ell}\left(1 + \frac{1}{E} - \frac{1}{E'}\right) \left[\frac{\mu(E)}{\mu(E')}\right]^n N_{n,\ell}^{\dagger}(E') = \\ &= \delta_{no} \delta_{\ell o} \mu_{ed}(E) + \int_{E_c}^{E_{i-1}} dE' \frac{E'}{E} K(E', E) P_{\ell}\left(1 + \frac{1}{E} - \frac{1}{E'}\right) \left[\frac{\mu(E)}{\mu(E')}\right]^n N_{n,\ell}^{\dagger}(E') + \\ &+ 2\delta_{\ell o} \frac{\mu_{pp}(E)}{E} \left[\frac{\mu(E)}{\mu(1)}\right]^n N_{n,o}^{\dagger}(1) + \omega_k \delta_{\ell o} \frac{\mu_{ph}(E)}{E} \left[\frac{\mu(E)}{\mu(E_1)}\right]^n N_{n,o}^{\dagger}(E_k) + \\ &+ \frac{(1-\delta_{no})}{2\ell+1} \mu(E) \left[(\ell+1) N_{n-1,\ell+1}^{\dagger}(E) + \ell N_{n-1,\ell-1}^{\dagger}(E) \right]. \end{aligned} \quad (6)$$

Since the right hand side of equation (6) does not contain $N_{n,\ell}^{\dagger}(E)$, we set it equal to some known quantity R . We can then rewrite Equation (6) in a more compact form,

$$\frac{R}{\mu(E)} = \int_{E_{i-1}}^E \left[\delta(E'-E) - \frac{E'}{E} \frac{K(E', E)}{\mu(E)} P_{\ell}\left(1 + \frac{1}{E} - \frac{1}{E'}\right) \left\{ \frac{\mu(E)}{\mu(E')} \right\}^n \right] N_{n,\ell}^{\dagger}(E') dE' \quad (7)$$

The integral is evaluated by writing the term in square brackets as a sum of two delta functions, namely

$$\begin{aligned} \delta(E'-E) - \frac{E'}{E} \frac{K(E', E)}{\mu(E)} P_{\ell}\left(1 + \frac{1}{E} - \frac{1}{E'}\right) \left[\frac{\mu(E)}{\mu(E')}\right]^n &= \alpha_{n\ell} \delta(E'-E_{i-1}) + \\ &+ \beta_{n\ell} \delta(E'-E_{n,\ell}^*) \end{aligned} \quad (8)$$

Equation (7) then becomes

$$\alpha_{nl} N_{n,l}^{\dagger}(E_{i-1}) + \beta_{n,l} N_{n,l}^{\dagger}(E_{n,l}^*) = \frac{R}{\mu(E)} \quad (9)$$

This procedure is somewhat different from that of Reference [7] in that a dependence on \underline{n} has been included. The three parameters $\alpha_{n,l}$, $\beta_{n,l}$, and $E_{n,l}^*$ may be determined by requiring that the first three moments of the scattering integral be given correctly,

$$I_{n,l}^j = \int_0^{E-E_{i-1}} (E-E')^j \frac{E'}{E} \frac{K(E',E)}{\mu(E)} p_l \left(1 + \frac{1}{E} - \frac{1}{E'}\right) \left[\frac{\mu(E)}{\mu(E')}\right]^n d(E-E') \quad (10)$$

Calculation of the first three moments of the functions on each side of equation (8) yields a system of equations:

$$\delta_{j0} - I_{n,l}^j = \alpha_{n,l} (E-E_{i-1})^j + \beta_{n,l} (E-E_{n,l}^*)^j, \quad j = 0, 1, 2 \quad (11)$$

After obtaining $\alpha_{n,l}$, $\beta_{n,l}$ and $E_{n,l}^*$, Equation (9) can be solved for $N_{n,l}^{\dagger}(E_{n,l}^*)$. Generally $E_{n,l}^*$ is not precisely equal to E and an interpolation must be performed to obtain the solution at E .

An interesting feature of Equation (5) has been pointed out by Morris in Reference [5] for the $n = l = 0$ case. If an energy deposition coefficient is defined as

$$\mu_{ed}(E) \equiv \mu(E) - \frac{2}{E} \mu_{pp}(E) - \frac{\omega_k}{E} \mu_{pH}(E) - \int_{E_c}^E dE' \frac{E'}{E} K(E',E), \quad (12)$$

then $N_{0,0}^{\dagger}(E) = 1.0$, for all E . This definition of $\mu_{ed}(E)$ is equivalent to the $\mu_K(E)$ defined by Hubbell [8]. By using this form for $\mu_{ed}(E)$ we

have an energy absorption coefficient which is consistent with the energy losses that are contained in the solution of the moments equations.

2. Description of Input

ADJMOM-I is written in FORTRAN-IV and is operational on the UNIVAC-1108, CDC-3600, and IBM-360/75 and 360/91. It requires 40K decimal words and two tape drives, one of which may be a scratch tape or disk. Several cases may be processed at one time with the output from each case being saved on the remaining tape.

The following are the input instructions for ADJMOM-I for the UNIVAC-1108 version. The only difference between the three versions of the code is in the assignment of the input and output units--which may be different at each installation.

3. Input

Card 1 (I5)*

NOUTI Number of energies at which adjoint moments will be punched out.

Card 2 (16F5.0)

(EOUTI(I), I = 1, NOUTI) Energy list for which adjoint moments will be punched out on cards. (MeV)

Card 3 (I1,15A4)

IPROB Problem type, IPROB = 1, Complete case; IPROB = 2, Use previous solution grid; IPROB = 3, Read new solution grid.

(TITLE(I), I = 1, 15) Title card for case.

* This is the format for the card or cards.

Card 4 (8I5)

NMUS Number of total cross sections.

NSORS Number of energies for energy deposition coefficient.

LS Maximum angular expansion coefficient used, $LS-1 = \ell_{\max}$.

MINNO Number of integration points for Equation (11).

MAXNO Number of integration points for scattering integral.

NPXS Number of partial cross sections. Pair production and each fluorescent X-ray are considered to be partial cross section reactions.

KEYIND Normally = 0, KEYIND = 1 causes the calculation of adjoint moments for the energy deposition coefficient in Equation (12) to be bypassed.

NEDI Number of energies for specifying the solution mesh.

Card 5 (3E10.0)

ETOP Maximum energy in the calculation. (MeV)

EBOTM Minimum energy in the calculation. (MeV)

XO Coefficient for converting input cross sections to units of TU/el. (Thomson units per electron)

Card 6 (8E10.0)

(EMUS(I), TABMU(I), I = 1, NMUS) EMUS(I) is the energy in MeV, TABMU(I) is the corresponding total cross section. TABMU usually has units of cm^2/g . The EMUS list must be in descending order.

Card 7 (8E10.0)

(ESORS(I), I = 1, NSORS) Energy list for tabulation of energy deposition coefficients (MeV).

Cards 8, 9, 10, 11, and 12 are required for NPXS > 0.

Card 8 (12I5)

(NPART(I), I = 1, NPXS) Number of energies for tabulating the
I-th partial cross section.

Card 9 (8E10.0)

(YIELD(I), I = 1, NPXS) Photon yield for each partial cross
section. Fluorescent yield for μ_{ph} and 2 for μ_{pp} .

Card 10 (8E10.0)

(ECUT(I), I = 1, NPXS) Lowest energy for each partial cross
section, MeV. 1.022 MeV for pair production and edge
energy for fluorescent gamma-rays.

Card 11 (8E10.0)

(ESTART(I), I = 1, NPXS) Energy for the gamma-ray produced by
I-th partial cross section interaction, MeV.

Card 12 (8E10.0) NPXS of these card sets are required. (I = 1, NPXS)

(EPART(J,I), SIGMA(J,I), J = 1, NPART(I)) EPART is the energy
in MeV and SIGMA is the corresponding partial cross
section. SIGMA has the same units as TABMU in Card 6.

Card 13 (4I5)

NLO Number of harmonic coefficients of the source angular
distribution. 1 for isotropic detector.

LZRO Index of source harmonic when NLO = 1. Otherwise
arbitrary.

MZRO Index of azimuthal harmonic. Usually MZRO = 0.

NNL Number of (n,l) combination for which solutions are
to be obtained.

Card 14 (4I3,4(I6,3I3))

(N(I), L(I), LINKH(I), LINKL(I), I = 1, NNL) where

N(I) n-index for the I-th moment.

L(I) ℓ -index for the I-th moment.

LINKH(I) I-index for the higher ℓ linkage moment.

LINKL(I) I-index for the lower ℓ linkage moment.

Card 15 (8E10.0)

(CNL(I), I = 1, NLO) Harmonic coefficients of the source angular distribution. For an isotropic source, CNL(1) = 1.0.

Card 16 (4(E10.0,I10))

(ED(I), NED(I), I = 1, NED1)

ED(I) Energy at which the grid of solution energies changes.

NED(I) Number of equally spaced intervals (logarithmically) between ED(I) and ED(I+1).

Cards 1 thru 15 are required for the calculation of adjoint moments for the energy deposition coefficient given by Equation (12). The following cards are required in order to do calculations for additional energy deposition coefficients.

Card 17 (15A4)

(TITLE(I), I = 1, 15) Title card for additional calculation.

Card 18 (I5,E10.0)

NSXXX Number of cross sections to be read. If NSXXX < 0, then card 19 must be supplied. Otherwise the cross sections for energy deposition will be assumed to be tabulated at the energy grid of the previous case.

DEN Changes cross sections to units of TU/el.

Card 19 (8E10.0) Not required for NSXXX > 0.

(ESORS(I), I = 1, ABS(NSXXX)) Energly list for tabulation of
energy deposition cross section, MeV.

Card 20 (8E10.0)

(STSORS(I), I = 1, ABS(NSXXX)) Energy deposition cross section
corresponding to the energy list, ESORS.

If NSXXX = 0, then the code assumes that the next card is a new
problem card, i.e. Card 3. When the calculation is completed with the
energy deposition coefficients given on Card 20, then the code expects
a title card (Card 15) and continues on with Card 16, etc. The code will
terminate when it encounters an end of file or the end of a data set,
whichever is appropriate for the machine being used.

4. Sample Input and Output

Input data for the calculation of adjoint moments for concrete
using μ_K for air taken from Reference [8].

| Card Type | CARD | COLUMNS |
|--------------|---|---------|
| | 11111111122222222233333333334444444445555555556666666667777777778 | |
| | 123456789012345678901234567890123456789012345678901234567890 | |
| 1 | 12 | |
| 2 | 1.0 0.80 .662 0.60 0.40 .279 0.20 0.10 0.08 0.06 0.04 0.02 | |
| 3 | 1 ADJOINT TEST PROBLEM--ALUMINUM--SKIP CONSISTENT PORTION | |
| 4 | 17 17 4 40 40 0 3 | |
| 5 | 1.0 0.01 5.181057 | |
| 6 | 1.0 0.0614 0.8 0.0683 0.6 0.0777 0.5 0.0841 | |
| 6 | 0.4 0.0922 0.3 0.1030 0.2 0.0120 0.15 0.1340 | |
| 6 | 0.1 0.162 0.08 0.189 0.06 0.255 0.05 0.334 | |
| 6 | 0.04 0.514 0.03 1.03 0.02 3.24 0.015 7.66 | |
| 6 | 0.01 25.8 | |
| 7 | 1.0 0.8 0.6 0.5 0.4 0.3 0.2 0.15 | |
| 7 | 0.1 0.08 0.06 0.05 0.04 0.03 0.02 0.015 | |
| 7 | 0.01 | |
| 13 | 1 0 0 10 | |
| 14 | 0 0 0 0 1 1 0 1 2 0 2 0 2 2 0 2 3 1 4 3 | |
| 14 | 4 0 5 0 3 3 0 4 4 2 7 5 5 1 8 6 6 0 9 0 | |
| 15 | 1.00 | |
| 16 | 1.0 45 0.10 70 0.01 | |
| 17 | ALUMINUM WITH MU-K (AIR) DETECTOR AS ADJOINT SOURCE | |
| 18 | 17 5.0012 | |
| 20 | 0.02797 0.02890 0.02958 0.02971 0.02952 0.02876 0.02677 0.02502 | |
| 20 | 0.02338 0.02427 0.03053 0.04062 0.06689 0.1480 0.51200 1.27100 | |
| 20 | 4.631 | |

IPR08= 1

ADJOINT TEST PROBLEM--ALUMINUM--SKIP CONSISTENT PORTION

| NWJS | NSORS | LS | INVO | MAXNO | NWJS | KEYINO | NE01 |
|------------------------------------|-----------|------------|------------|---------|-----------|------------|------------|
| 17 | 17 | 4 | 40 | 40 | 0 | 1 | 3 |
| 1.00000 | | .01000 | | 5.18106 | | | |
| TOTAL CROSS SECTIONS ENERGY--SIGMA | | | | | | | |
| .100+01 | .614-01 | .900+00 | .693-01 | .600+00 | .777-01 | .500+00 | .841-01 |
| .300+00 | .103+00 | .200+00 | .120+00 | .150+00 | .134+00 | .100+00 | .162+00 |
| .600-01 | .255+00 | .500-01 | .334+00 | .400-01 | .514+00 | .300-01 | .103+01 |
| .150-01 | .766+01 | .100-01 | .258+02 | | | | |
| TOTAL CROSS SECTIONS ENERGY--SIGMA | | | | | | | |
| .196+01 | .318+00 | .157+01 | .354+00 | .117+01 | .403+00 | .978+00 | .436+00 |
| .587+00 | .534+00 | .391+00 | .622+00 | .294+00 | .694+00 | .196+00 | .839+00 |
| .117+00 | .132+01 | .978-01 | .173+01 | .783-01 | .266+01 | .587-01 | .534+01 |
| .294-01 | .397+02 | .196-01 | .134+03 | | | | |
| ENERGIES FOR ENTERING DO E DATA | | | | | | | |
| .100+01 | .800+00 | .00+00 | .500+00 | .400+00 | .300+00 | .200+00 | .150+00 |
| .600-01 | .500-01 | .400-01 | .300-01 | .200-01 | .150-01 | .100-01 | |
| 1 | 0 | 0 | 0 | 0 | .00000 | .00000 | .10000+01 |
| 2 | 0 | 1 | 1 | 1 | .00000 | .33333+00 | .33333+00 |
| 3 | 2 | 0 | 2 | 0 | .00000 | .00000 | .33333+00 |
| 4 | 0 | 2 | 2 | 2 | .00000 | .00000+00 | .13333+00 |
| 5 | 4 | 3 | 3 | 1 | .66667+00 | .33333+00 | .20000+00 |
| 6 | 5 | 0 | 4 | 0 | .10000+01 | .00000 | .20000+00 |
| 7 | 0 | 4 | 3 | 3 | .00000 | .42857+00 | .57143-01 |
| 8 | 7 | 5 | 4 | 2 | .60000+00 | .40000+00 | .11429+00 |
| 9 | 8 | 6 | 5 | 1 | .66667+00 | .33333+00 | .14286+00 |
| 10 | 9 | 0 | 6 | 0 | .10000+01 | .00000 | .14286+00 |
| .1000+01 45 | | | | | | | |
| .1000+00 70 | | | | | | | |
| .1000-01 0 | | | | | | | |
| I | ENERGY(I) | TOTAL(I) | WU-ED(I) | I | ENERGY(I) | TOTAL(I) | WU-ED(I) |
| | (MC*2) | (TU/EL) | (TU/EL) | | (VC*2) | (TU/EL) | (TU/EL) |
| 1 | 1.956926 | .3181 7 | .140335 | 2 | 1.859311 | .325981 | .141793 |
| 4 | 1.678447 | .342298 | .143528 | 5 | 1.594724 | .350761 | .144612 |
| 7 | 1.439597 | .367422 | .145940 | 8 | 1.367788 | .375946 | .146507 |
| 10 | 1.234737 | .393592 | .147659 | 11 | 1.173146 | .402718 | .148247 |
| 13 | 1.059028 | .421014 | .148890 | 14 | 1.006203 | .430471 | .149235 |
| 16 | .908324 | .449289 | .149264 | 17 | .863016 | .458863 | .149166 |
| 19 | .779066 | .478567 | .149986 | 20 | .740205 | .488089 | .148704 |
| 22 | .668202 | .5077 5 | .147009 | 23 | .634471 | .517807 | .146615 |
| 25 | .573114 | .5385 1 | .145281 | 26 | .548526 | .548993 | .144635 |
| 28 | .491557 | .5705 4 | .143584 | 29 | .467038 | .581680 | .143205 |
| 31 | .421607 | .604546 | .142814 | 32 | .400576 | .616313 | .142833 |
| 34 | .361610 | .640886 | .143729 | 35 | .343573 | .653589 | .144530 |
| 37 | .310152 | .679755 | .146832 | 38 | .294681 | .693228 | .148371 |
| 40 | .266016 | .726999 | .158342 | 41 | .252747 | .744618 | .164125 |
| 43 | .228161 | .781149 | .177178 | 44 | .216740 | .800080 | .184497 |
| 46 | .195693 | .839331 | .200846 | 47 | .189306 | .858622 | .212881 |
| 49 | .177303 | .898585 | .238462 | 50 | .171566 | .919197 | .252836 |
| 52 | .160642 | .961936 | .280818 | 53 | .155444 | .986501 | .298512 |
| 55 | .145547 | 1.056439 | .354926 | 56 | .140437 | 1.093245 | .385090 |
| 58 | .131870 | 1.170751 | .449550 | 59 | .127603 | 1.211540 | .443945 |
| 61 | .119478 | 1.297432 | .557305 | 62 | .115612 | 1.351792 | .605532 |
| 64 | .104251 | 1.490055 | .731805 | 65 | .104748 | 1.564402 | .800299 |
| 67 | .098079 | 1.724410 | .948880 | 68 | .094905 | 1.835581 | 1.054479 |
| 70 | .088862 | 2.084345 | 1.292397 | 71 | .085087 | 2.221097 | 1.423865 |
| 73 | .080512 | 2.522108 | 1.714607 | 74 | .077907 | 2.693764 | 1.841274 |
| 76 | .072946 | 3.157862 | 2.335696 | 77 | .070586 | 3.419086 | 2.592236 |
| 79 | .066091 | 4.008147 | 3.172193 | 80 | .063953 | 4.330707 | 3.499329 |
| 82 | .059881 | 5.087379 | 4.234472 | 83 | .057943 | 5.537917 | 4.644091 |
| 85 | .054254 | 6.466963 | 5.408616 | 86 | .052498 | 7.194455 | 6.454608 |
| 88 | .049156 | 8.152229 | 7.442965 | 89 | .047565 | 9.674111 | 8.798207 |
| 91 | .044537 | 11.651073 | 10.768286 | 92 | .043095 | 12.786255 | 11.900105 |
| 94 | .040351 | 15.399203 | 14.506628 | 95 | .039046 | 16.906170 | 14.010516 |
| 97 | .036560 | 20.542615 | 19.641100 | 98 | .035377 | 22.710574 | 21.806987 |
| 100 | .033124 | 27.649263 | 26.730476 | 101 | .032052 | 30.507823 | 29.595175 |
| 103 | .030012 | 37.142109 | 36.224645 | 104 | .029040 | 40.243886 | 40.063935 |
| 106 | .027191 | 49.909615 | 48.985024 | 107 | .026311 | 55.076927 | 54.150141 |
| 109 | .024636 | 67.071924 | 66.140837 | 110 | .023439 | 74.016123 | 73.082260 |
| 112 | .022321 | 90.135813 | 89.198835 | 113 | .021599 | 99.467899 | 98.528371 |
| 115 | .020224 | 121.130632 | 120.385306 | 116 | .019569 | 133.671711 | 133.671711 |

ALUMINUM WITH MU-K (AIR) DETECTOR AS ADJOINT SOURCE

17 5.0012000+00

ENERGIES FOR NEW SOURCE FUNCTIONS

| | | | | | | | |
|---------|---------|---------|---------|---------|---------|---------|---------|
| .196+01 | .157+01 | .117+01 | .978+00 | .783+00 | .587+00 | .391+00 | .294+00 |
| .196+00 | .157+00 | .117+00 | .978-01 | .783-01 | .587-01 | .391-01 | .294-01 |
| .196-01 | | | | | | | |

DOSE COEFFICIENTS FOR ADJOINT SOURCE

| | | | | | | | |
|---------|---------|---------|---------|---------|---------|---------|---------|
| .280-01 | .289-01 | .296-01 | .297-01 | .295-01 | .289-01 | .268-01 | .250-01 |
| .234-01 | .243-01 | .305-01 | .406-01 | .669-01 | .149+00 | .512+00 | .127+01 |
| .463+01 | | | | | | | |

| I | ENERGY(I) (MC*2) | TOTAL(I) (TU/EL) | MU-ED(I) (TU/EL) | I | ENERGY(I) (MC*2) | TOTAL(I) (TU/EL) | MU-ED(I) (TU/EL) | I | ENERGY(I) (MC*2) | TOTAL(I) (TU/EL) | MU-ED(I) (TU/EL) |
|-----|---------------------|---------------------|---------------------|-----|---------------------|---------------------|---------------------|-----|---------------------|---------------------|---------------------|
| 1 | 1.956926 | .318117 | .139884 | 2 | 1.859311 | .325981 | .140937 | 3 | 1.766566 | .334040 | .14199A |
| 4 | 1.678447 | .342298 | .143067 | 5 | 1.594724 | .350761 | .144144 | 6 | 1.515176 | .359091 | .146917 |
| 7 | 1.439597 | .367422 | .145518 | 8 | 1.367788 | .375946 | .146121 | 9 | 1.299561 | .384668 | .146727 |
| 10 | 1.234737 | .393352 | .147335 | 11 | 1.173146 | .402718 | .147039 | 12 | 1.11462A | .411764 | .14A121 |
| 13 | 1.059028 | .421014 | .148303 | 14 | 1.006203 | .430471 | .148486 | 15 | .956012 | .439915 | .14A4A7 |
| 16 | .908324 | .442889 | .148268 | 17 | .863016 | .458863 | .148050 | 18 | .819967 | .468642 | .147833 |
| 19 | .779066 | .478567 | .147572 | 20 | .740205 | .48A089 | .146A89 | 21 | .70328A | .497801 | .146209 |
| 22 | .668202 | .507705 | .145532 | 23 | .63A871 | .517807 | .144A59 | 24 | .603203 | .52A110 | .14418A |
| 25 | .573114 | .538511 | .143223 | 26 | .544526 | .548993 | .141933 | 27 | .517364 | .559679 | .140655 |
| 28 | .491957 | .570574 | .139388 | 29 | .467038 | .581680 | .138138 | 30 | .443741 | .593003 | .13688A |
| 31 | .421607 | .604546 | .135655 | 32 | .400576 | .616313 | .134433 | 33 | .380595 | .628430 | .133005 |
| 34 | .361610 | .640886 | .131416 | 35 | .343573 | .653589 | .129A45 | 36 | .326435 | .666543 | .128293 |
| 37 | .310152 | .679755 | .126759 | 38 | .2946A1 | .693228 | .125244 | 39 | .279982 | .709797 | .124145 |
| 40 | .266016 | .726999 | .123087 | 41 | .252747 | .744618 | .122038 | 42 | .240139 | .762665 | .120999 |
| 43 | .228161 | .781149 | .119968 | 44 | .216780 | .A00080 | .118946 | 45 | .205966 | .A19471 | .117933 |
| 46 | .195693 | .839331 | .115928 | 47 | .189360 | .A58622 | .117574 | 48 | .183233 | .878357 | .11A223 |
| 49 | .177303 | .898545 | .118876 | 50 | .171566 | .919197 | .119532 | 51 | .166014 | .940324 | .120193 |
| 52 | .160642 | .961936 | .120855 | 53 | .155444 | .986501 | .122070 | 54 | .150414 | 1.020871 | .125315 |
| 55 | .145547 | 1.056439 | .128647 | 56 | .140837 | 1.093245 | .132067 | 57 | .136280 | 1.131335 | .13557A |
| 58 | .131870 | 1.170751 | .139182 | 59 | .127603 | 1.211540 | .142A82 | 60 | .123474 | 1.253751 | .146681 |
| 61 | .119478 | 1.297432 | .150580 | 62 | .115612 | 1.331792 | .142A82 | 63 | .111871 | 1.419241 | .164704 |
| 64 | .108251 | 1.490055 | .173412 | 65 | .104748 | 1.564402 | .182580 | 66 | .10135A | 1.642459 | .192233 |
| 67 | .098079 | 1.724410 | .202396 | 68 | .094905 | 1.835581 | .217493 | 69 | .091834 | 1.956012 | .234087 |
| 70 | .088862 | 2.084345 | .251948 | 71 | .085947 | 2.221097 | .271171 | 72 | .083204 | 2.366822 | .291861 |
| 73 | .080512 | 2.522108 | .314129 | 74 | .077907 | 2.693764 | .338940 | 75 | .075386 | 2.916596 | .371159 |
| 76 | .072946 | 3.157862 | .406440 | 77 | .070586 | 3.419086 | .445074 | 78 | .06A302 | 3.70191A | .4A7382 |
| 79 | .066091 | 4.008147 | .533710 | 80 | .063953 | 4.339707 | .594443 | 81 | .061893 | 4.699695 | .63999A |
| 82 | .059881 | 5.008379 | .700834 | 83 | .057943 | 5.537917 | .770481 | 84 | .05606A | 6.077485 | .A52099 |
| 85 | .054254 | 6.669623 | .942363 | 86 | .052498 | 7.319455 | 1.042188 | 87 | .050800 | 8.032601 | 1.15258A |
| 88 | .049156 | 8.815229 | 1.274683 | 89 | .047565 | 9.674111 | 1.409711 | 90 | .046026 | 10.616675 | 1.559043 |
| 91 | .044537 | 11.651073 | 1.724194 | 92 | .043095 | 12.786255 | 1.906A39 | 93 | .041701 | 14.032040 | 2.10A833 |
| 94 | .040351 | 15.399203 | 2.332224 | 95 | .039046 | 16.906170 | 2.579A88 | 96 | .0377A2 | 18.654040 | 2.862540 |
| 97 | .036560 | 20.582615 | 3.176160 | 98 | .035377 | 22.710578 | 3.524140 | 99 | .034232 | 25.05A54A | 3.910245 |
| 100 | .033124 | 27.649263 | 4.338651 | 101 | .032052 | 30.507823 | 4.813993 | 102 | .031015 | 33.661923 | 5.341414 |
| 103 | .030012 | 37.142109 | 5.926620 | 104 | .029040 | 40.983886 | 6.577038 | 105 | .028101 | 45.257091 | 7.805414 |
| 106 | .027191 | 49.909615 | 8.113345 | 107 | .026311 | 55.076927 | 9.010626 | 108 | .025460 | 60.779237 | 10.007142 |
| 109 | .024636 | 67.071924 | 11.113864 | 110 | .023839 | 74.016123 | 12.342084 | 111 | .02306A | 81.679272 | 13.708036 |
| 112 | .022331 | 90.135813 | 15.224053 | 113 | .021599 | 99.467899 | 16.907733 | 114 | .020900 | 109.766159 | 1A.777615 |
| 115 | .020224 | 121.130632 | 20.854295 | 116 | .019569 | 133.671711 | 23.160440 | | | | |

[illegible][illegible][illegible]15

INTERPOLATED MOMENTS AT MONOENERGETIC SOURCES
ALUMINUM WITH MU-K (AIR) DETECTOR AS ADJOINT SOURCE

| | | | | | | | | | |
|---------|-------------|-----------|-------------|--------|-------------|-------|-------------|---------|--------------|
| E(MEV)= | 1.000 | E(MC**2)= | 1.957 | TOTAL= | .318117+00 | D05E= | .139884+00 | DIRECT= | .439724 |
| 0 0 | .9301593+00 | 1 1 | .4464763+00 | 2 0 | .7454428+00 | 2 2 | .2203788+00 | 3 1 | .5376884+00 |
| 4 0 | .7844774+00 | 3 3 | .1072226+00 | 4 2 | .3388688+00 | 5 1 | .6322624+00 | 6 0 | .8507860+00 |
| E(MEV)= | .800 | E(MC**2)= | 1.566 | TOTAL= | .353745+00 | D05E= | .144423+00 | DIRECT= | .408269 |
| 0 0 | .9134349+00 | 1 1 | .4397546+00 | 2 0 | .7732403+00 | 2 2 | .2156611+00 | 3 1 | .5527852+00 |
| 4 0 | .8433346+00 | 3 3 | .1038927+00 | 4 2 | .3433221+00 | 5 1 | .6729090+00 | 6 0 | .9410669+00 |
| E(MEV)= | .662 | E(MC**2)= | 1.295 | TOTAL= | .385210+00 | D05E= | .146764+00 | DIRECT= | .380999 |
| 0 0 | .8962830+00 | 1 1 | .4310800+00 | 2 0 | .7923109+00 | 2 2 | .2097657+00 | 3 1 | .5599288+00 |
| 4 0 | .8876797+00 | 3 3 | .1000873+00 | 4 2 | .3426243+00 | 5 1 | .6996735+00 | 6 0 | .1011777+01 |
| E(MEV)= | .600 | E(MC**2)= | 1.174 | TOTAL= | .402563+00 | D05E= | .147928+00 | DIRECT= | .367466 |
| 0 0 | .8860015+00 | 1 1 | .4253234+00 | 2 0 | .8003863+00 | 2 2 | .2060183+00 | 3 1 | .5617157+00 |
| 4 0 | .9095753+00 | 3 3 | .9778213-01 | 4 2 | .3409076+00 | 5 1 | .7118044+00 | 6 0 | .1048960+01 |
| E(MEV)= | .400 | E(MC**2)= | .783 | TOTAL= | .477638+00 | D05E= | .147596+00 | DIRECT= | .309013 |
| 0 0 | .8290155+00 | 1 1 | .3916825+00 | 2 0 | .8101221+00 | 2 2 | .1857453+00 | 3 1 | .5483321+00 |
| 4 0 | .9667153+00 | 3 3 | .8625194-01 | 4 2 | .3211525+00 | 5 1 | .7296182+00 | 6 0 | .1161190+01 |
| E(MEV)= | .279 | E(MC**2)= | .546 | TOTAL= | .548441+00 | D05E= | .142000+00 | DIRECT= | .258916 |
| 0 0 | .7574802+00 | 1 1 | .3482946+00 | 2 0 | .7688676+00 | 2 2 | .1616742+00 | 3 1 | .4995899+00 |
| 4 0 | .9331244+00 | 3 3 | .7376906-01 | 4 2 | .2837232+00 | 5 1 | .6773000+00 | 6 0 | .1136241+01 |
| E(MEV)= | .200 | E(MC**2)= | .391 | TOTAL= | .621781+00 | D05E= | .133783+00 | DIRECT= | .215162 |
| 0 0 | .6652552+00 | 1 1 | .2956628+00 | 2 0 | .6740279+00 | 2 2 | .1346295+00 | 3 1 | .4212863+00 |
| 4 0 | .8091902+00 | 3 3 | .6054551-01 | 4 2 | .2337241+00 | 5 1 | .5670974+00 | 6 0 | .9766046+00 |
| E(MEV)= | .100 | E(MC**2)= | .196 | TOTAL= | .839331+00 | D05E= | .116928+00 | DIRECT= | .139311 |
| 0 0 | .4080004+00 | 1 1 | .1655111+00 | 2 0 | .3408293+00 | 2 2 | .7251487-01 | 3 1 | .2010388+00 |
| 4 0 | .3503448+00 | 3 3 | .3200693-01 | 4 2 | .1092809+00 | 5 1 | .2362071+00 | 6 0 | .37233348+00 |
| E(MEV)= | .080 | E(MC**2)= | .157 | TOTAL= | .981135+00 | D05E= | .121806+00 | DIRECT= | .124149 |
| 0 0 | .3233917+00 | 1 1 | .1260858+00 | 2 0 | .2355980+00 | 2 2 | .5460314-01 | 3 1 | .1373500+00 |
| 4 0 | .2200268+00 | 3 3 | .2396956-01 | 4 2 | .7516022-01 | 5 1 | .1481498+00 | 6 0 | .2165608+00 |
| E(MEV)= | .060 | E(MC**2)= | .117 | TOTAL= | .132593+01 | D05E= | .153651+00 | DIRECT= | .115882 |
| 0 0 | .2389711+00 | 1 1 | .8842998-01 | 2 0 | .1420861+00 | 2 2 | .3747006-01 | 3 1 | .8222843-01 |
| 4 0 | .1172337+00 | 3 3 | .1632625-01 | 4 2 | .4525368-01 | 5 1 | .7921132-01 | 6 0 | .1051857+00 |
| E(MEV)= | .040 | E(MC**2)= | .078 | TOTAL= | .266830+01 | D05E= | .335244+00 | DIRECT= | .125639 |
| 0 0 | .1784045+00 | 1 1 | .6207335-01 | 2 0 | .7882836-01 | 2 2 | .2554569-01 | 3 1 | .4592844-01 |
| 4 0 | .5490350-01 | 3 3 | .1102768-01 | 4 2 | .2570492-01 | 5 1 | .3777836-01 | 6 0 | .4351054-01 |
| E(MEV)= | .020 | E(MC**2)= | .039 | TOTAL= | .167927+02 | D05E= | .256118+01 | DIRECT= | .152517 |
| 0 0 | .1611192+00 | 1 1 | .5393671-01 | 2 0 | .5637190-01 | 2 2 | .2167714-01 | 3 1 | .3347566-01 |
| 4 0 | .3471667-01 | 3 3 | .9297572-02 | 4 2 | .1905338-01 | 5 1 | .2048446-01 | 6 0 | .2524607-01 |

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| | | | | 14. Sponsoring Agency Code |
| 15. SUPPLEMENTARY NOTES | | | | |
| 16. ABSTRACT (A 200-word or less factual summary of most significant information. If document includes a significant bibliography or literature survey, mention it here.) In this paper we discuss a computer code for generating spatial-angular moments of the adjoint gamma-ray flux in an infinite medium. The equation for the flux moments is given and techniques used for the solution are described. Details of the input data and a sample problem are also supplied. | | | | |
| 17. KEY WORDS (Alphabetical order, separated by semicolons) Adjoint; buildup factor; dosimetry; gamma-ray transport; moment methods; shielding. | | | | |
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